(almost) All Life occurs in a Plasma of Spherical Ions in water Na⁺, K⁺, Ca⁺⁺, and Cl⁻

each with a different diameter

Ion Diameters Pauling Diameters			
Ca ⁺⁺	1.98 Å		
Na⁺	2.00 Å		
K+	2.66 Å		

lons are involved in most of biology

lons are controlled by ion channels that are natural nano-valves*

lons control all electrical activity in cells

lons produce signals of the nervous system

lons coordinate contraction in skeletal muscle

lons coordinate contraction in the heart, allowing the heart to function as a pump

lons are involved in secretion and absorption in all cells: kidney, intestine, liver, adrenal glands, etc.

lon channels are involved in thousands of diseases and many drugs act on channels

lon channels are proteins whose genes (blueprints) can be manipulated by molecular genetics

lon channels have structures shown by x-ray crystallography in favorable cases.



*nearly <u>pico</u>-valves: diameter is 400 - 900 pico-meters

Charged Particles in a Dielectric with Friction is THE Fundamental Problem in Plasmas, including Plasmas of Life

I am not qualified to discuss the importance of this problem in physical plasmas but to an outsider, it seems fundamental.

Seeking a Simple Description of the Brownian Motion of <u>Charged</u> Particles

Einstein's Brownian Particles are Uncharged

(nearly) Everything Dissolved in Water is Charged (somewhere)

Conjecture Fluctuations in charge density are a significant –even dominant– source of Fluctuations in Plasmas

but

Einstein's treatment of Brownian motion does not discuss charge

Simplified Descriptions are Clearly Possible:

Ohm's Law works well for a wide range of ionic solutions (Ohm's law involves ONLY charge)

Fick's law works well for a wide range of nonionic solutions (Fick's law involves only mass)

Self-consistent Simulation

Consider a random process in which charged particles move in an electric field created by their own charge and charge applied by boundary conditions. (No other applied fields are allowed)

- 1) Start with an overall neutral system
 - 2) Choose a small volume
 - 3) Count the number of particles of each type in that volume
 - 4) Compute the electric field from the location and amount of charges
 - 5) Allow the particles to move a small amount
 - 6) Count again, etc.

Construct graphs of number density ('concentration') of particles vs. time and location.

Typical Time Series

[X] = number of X	Time =	1	2	3	•••
Number of Na ⁺	[Na⁺]	7	6	[Na ⁺]	
Number of K ⁺	[K+]	3	2	[K+]	
Number of CI [−]	[CI ⁻]	9	9	[CI ⁻]	•••
Number of Positive Charg	jes	10	8	[Na ⁺]+[K ⁺]	
Number of Negative Chai	rges	9	9	[CI ⁻]	
Net Charge Q (units: number of cl	harges)	+1	-1	[Na⁺]+[K⁺]-[CI [−]]	
Number of Particles N		19	17	[Na⁺]+[K⁺]+[CI [−]]	

Typical Time Series

	Time =	1	2	3	•••
1) Number of Na ⁺	[Na⁺]	0	6	[Na⁺]	
2) Number of K ⁺	[K+]	3	2	[K+]	
3) Number of CI [−]	[CI ⁻]	9	9	[CI ⁻]	

Gives equation for [K⁺]
 Gives equation for [Na+]
 Gives equation for [Cl⁻]

Variables [Na⁺], [K⁺], [Cl⁻] are highly correlated

so we have severe 'closure' problems

Time Series of

Time =	1	2	3	•••
1) Net Charge Q (units: number of charges)	+1	-1	[Na ⁺]+[K ⁺]-[Cl ⁻]	
2) Number of Particles N	19	17	[Na⁺]+[K⁺]+[Cl ⁻]	•••

Gives equation for *Q* Gives equation for *N*

Variables

 $Q = [Na^+]+[K^+]-[CI^-]$ $N = [Na^+]+[K^+]+[CI^-]$

are almost uncorrelated

(we know from experiments and common sense)

so (I imagine) we have almost no closure problems



We know PDE's for $[Na^+]$, $[K^+]$, $[CI^-]$.

What are the PDE's for *Q*, *N*, and [Cl⁻]?



How do we "change variable"?

How do we construct the counting process for charge and density?

Charged Brownian Motion in Langevin Form



Similar Equation for location $x(t)_{k}^{-}$ of negative species k



What has been done?

We start with Langevin equations of charged particles

Simplest stochastic trajectories are Brownian Motion of Charged Particles

Gouy-Chapman, (nonlinear) Poisson-Boltzmann, Debye-Hückel, are models with similar resolution but constrained to equilibrium, i.e., zero flux of all species. **Devices do not exist at equilibrium**

Once we learn to count Trajectories of Brownian Motion, we can count trajectories of <u>Molecular Dynamics</u>

Equilibrium

Configurations Boltzmann Distribution $\lim N, V \rightarrow \infty$



<u>Nonequilibrium</u>

Trajectories Fokker Planck Equation Finite OPEN System

Schuss, Nadler, and Eisenberg

Langevin Equations



Electric Force from Poisson Equation



From Trajectories to Probabilities Main Result: Theory of Stochastic Processes

Joint probability density of position and velocity

 $p(\tilde{x}, \tilde{v}) = \Pr\left\{\left\{x_{j}, v_{j}\right\}_{j=1}^{2N}\right\}; \quad N = \text{Number of Particles}$ satisfies a Fokker Planck equation

$$0 = \sum_{j} \mathsf{L}_{j}^{p} p(\tilde{x}, \tilde{v}) + \sum_{j} \mathsf{L}_{j}^{n} p(\tilde{x}, \tilde{v})$$

with Fokker Planck Operator

$$\mathsf{L}_{j}^{c} p = -v_{j}^{c} \cdot \nabla_{x_{j}^{c}} p + \nabla v_{j}^{c} \cdot \left(\gamma v_{j}^{c} - \frac{f_{j}^{c}}{m_{j}^{c}}\right) p + \nabla \cdot \nabla_{v_{j}^{c}} \frac{\gamma kT}{m_{j}^{c}} p$$

Coordinates are positions and velocities of N particles in 12N dimensional phase space

Schuss, Nadler, and Eisenberg

Conditional PNP

Electric Force $abla \phi$ depends on Conditional Density of Charge

$$\nabla_{y} \cdot \begin{bmatrix} \varepsilon_{0} \varepsilon(y) \\ e \end{bmatrix} \nabla_{y} \phi(y | x) = P(y)$$
Closures or Approximations
Needed
Needed
$$(p | x) = P(y)$$
Channel
Protein
$$(p | x) - \rho_{-}(y | x)$$

Nernst-Planck gives UNconditional Density of Charge

$$\nabla_{x} \cdot \left[\frac{1}{m\gamma(x)} \rho_{+}^{*}(x) \left[e \nabla_{y} \overline{\phi}(y | x) \Big|_{y=x} - DBF \right] \right] = 0$$
Mass
Friction

Schuss, Nadler, Eisenberg



Counting at low resolution gives <u>'Semiconductor Equations'</u>

Poisson-Nernst-Planck (PNP)

Only contains correlations of means

Gouy-Chapman, (nonlinear) Poisson-Boltzmann, Debye-Hückel, are siblings with similar resolution but without current or flux of any species Devices do not exist at equilibrium



Poisson's Equation

$$-\varepsilon_0 \nabla \cdot \left(\varepsilon(\mathbf{x}) \nabla \phi(\mathbf{x}) \right) = e \mathbf{P}(\mathbf{x}) + e \sum_i z_i \rho_i(\mathbf{x})$$

Channel Protein

Drift-diffusion & Continuity Equation

$$\nabla \cdot \mathbf{J}(\mathbf{x}) = 0; \quad -\mathbf{J}_i(\mathbf{x}) = D_i(\mathbf{x})\rho_i(\mathbf{x})\frac{1}{kT}\nabla\mu_i(\mathbf{x})$$

Chemical Potential



$$\mu_i(\mathbf{x}) = z_i e\phi(\mathbf{x}) + kT \ln\left(\frac{\rho_i(\mathbf{x})}{\rho^*}\right) + \mu_i^{ex}(\mathbf{x})$$

Special Chemistry

Solving semiconductor equations requires a trick



Or much better (but much harder) Newton Iteration

Electrodiffusion of charged, hard spheres

Correlations put in by Hand



How well can we do biology with correlations done by hand?

