**Mass Action**and **Conservation of Charge**

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**Abstract**

Electronic technology has remade our world, in fifty years, increasing capability by factors of billions, unprecedented in human history. Electronics depends on models that precisely represent its devices and circuit. Circuit models are robust and work well over a wide range of conditions without changing parameters.

Circuit models are based on the laws of electricity, conservation of charge, and Kirchoff’s current law: current flow is exactly the same everywhere in a sequence of devices (with two terminals, like diodes or ionic channels).

Chemical and biochemical models have been built on conservation of matter, expressed as the law of mass action, with constant rate constants. Chemical and biochemical technology has made great progress but not by making devices, rather by making compounds.

Device design in the chemical world is difficult because the law of mass action is rarely transferrable from one set of conditions to another. It is not robust in most cases and cannot serve the same role that circuit models do in our electrical technology.

The law of mass action does not conserve charge, as is shown mathematically and discussed at length in the paper, along with the abstract nature of charge. The mathematical methods to combine the laws of mass action and conservation of charge have not been available for dissipative systems like ions in water.

I believe robust models and device design in the chemical world will not be possible until the law of mass action is extended to conserve charge. Variational methods have only recently been extended to include dissipation (by combining two variations). The Energy Variational Approach ***EnVarA*** developed by Chun Liu, more than anyone else, should allow the development of robust models of chemical, biochemical, and biological systems.

Scientific laws are sometimes vague, sometimes precise. Scientific laws are often residues of revolutions that once gave us new knowledge of the world around and within us. Scientific laws are usually learned early in our careers before we have refereed grants and papers, before our critical skills were practiced. It is easy to continue to use scientific laws uncritically even after they have been overtaken by new knowledge.

**Conservation of Charge and Conservation of Mass.** I write to discuss two scientific laws: conservation of charge, and conservation of mass. Both seem precise laws of science and in historical context they certainly once were precise laws, not vague residues of previous revolutions. Both laws are so widely used they might be called foundations of their fields: chemistry and physics.

Chemistry uses mass conservation almost everywhere, in the form of the law of mass action to describe chemical reactions and binding. Physics uses conservation of charge almost everywhere it deals with electricity. For example, Kirchoff’s current law says the flow of charge is continuous, without loss in circuits. Electrical engineering makes continual use of Kirchoff’s laws. Kirchoff’s laws allow the robust calculation of electronic devices that behave in one way, with one set of parameters, over a range of conditions. Without Kirchoff’s laws, we would not have integrated circuits or digital devices. We would not have computers, smartphones, and the video devices that have remade our world.

**Biological systems involve both chemical reactions and charge** and so must deal with both the law of mass action and conservation of charge. Biological systems are always embedded in ionic solutions, and nearly always involve chemical reactants and enzymes with electrical charge. Substrates of reactions catalyzed by enzymes are usually charged and are always embedded in solutions containing the ‘bio-ions’ Na+, K+, and Cl− and involving Ca2+, often as a variable that turns biological function on or off. But we face a problem when we try to apply both conservation laws and mass action laws together. The laws conflict.

**My goal** is to replicate the success of electronic models in chemistry, biochemistry, and biology. The goal is to produce robust theories (or simulations) of chemical reactions in ionic solutions that successfully use one set of parameters in a range of conditions, and include the global properties of the electric field. My goal is to produce electrochemical devices and technology as robust as those of electronics. I believe that goal can only be reached after the law of mass action is extended to conserve charge.

**The Law of Mass Action** has been a keystone of classical chemistry. It was developed to deal with isolated systems, originally with perfect gases in mind [3, 30, 54, 63, 65], and allowed chemists to deal with molecular and atomic reactions before physicists were convinced that atoms existed. It allows chemists to build molecules of great use and importance.

But **isolated systems must make contact** with the outside world to interact with it just as electronic systems make contact with inputs and outputs and biological systems make contact with surrounding solutions and cells. Interactions with the outside world almost always involve electricity because charge flows through the contacts that connect chemical reactions with the outside world. **The law of mass action was not designed to contact the outside world.** It was designed to help chemists build and understand individual molecules.

Engineering deals with systems that are not isolated. Its device contact the world through power supplies. Devices have outputs that follow inputs according to a simple rule, for example, the gain of an amplifier. Kirchoff’s current law and conservation of charge are the bases of those rules.

Biology deals with systems that are not isolated. They usually have inputs and outputs and are driven by concentration gradients that are power supplies. Biological systems interact with surrounding solutions, cells, and tissues.

I argue that the law of mass action must be extended to deal with inputs and outputs and electrical charge if theories and simulations of nanodevices (technological or biological) are to be useful in more than one set of conditions.

**The Law of Mass Action does not conserve charge.** Conservation of matter does not imply conservation of charge because (displacement) charge can move in a vacuum and one atom of one mass can have different charges. Displacement current is “Maxwell’s extension of Ampere’s law” (p. 521 of [41]) needed to make a successful theory of electromagnetism, including the propagation of light. [22, 40, 41, 74]

In my view, the law of mass action must be extended before it can deal with the ions and charges that are almost always present in biological systems. If the law of mass action does not conserve charge it cannot be used to derive robust models that describe charge movement. Charge movement at power supplies, inputs and outputs are always involved as devices contact the outside world. Charge must be included if a chemical reaction connects to the outside world through electrical inputs and outputs like those used in our engineering technology. Charge has not been included in chemical reactions in my opinion because the mathematics needed was not available until recently. Mathematics that forces chemical reactions to conserve charge in ionic solutions has been developed only in the last few years as variational principles have been extended (by taking a second variation) to include dissipative systems like ionic solutions, by Chun Liu, more than anyone else.

That new mathematics uses variational principles to derive field equations—usually coupled partial differential equations—that automatically and universally satisfy conservation of mass and conservation of charge. The variational principle is defined by the Euler-Lagrange process [21]. Only recently has that variational principle been generalized by mathematicians into an energy variational approach ***EnVarA*** that combines two variations and thus can deal with dissipation and ionic solutions. Ionic solutions experience friction because an ionic solution is a condensed phase essentially without empty space. ***EnVarA*** deals with conserved energy described by the classical Hamiltonians described in textbooks of mechanics. ***EnVarA*** deals with friction. It allows energy to be degraded into entropy and describes that process by a combination of the classical Hamiltonian of conservative systems and Rayleigh’s dissipation of frictional systems.

**How has the law of mass action produced such useful chemistry, if it does not conserve charge?** The law of mass action has been immensely productive chemistry as a guide to understand and build molecules, not as a foundation to build devices.

The law of mass action has been productive in biochemistry even though it does not conserve charge because it is used in a such a clever and ingenious way. A senior biochemist recently told me “The art of biochemistry is to study enzymes in conditions that give insight to biological function, even if the conditions are not just those inside cells. The art is to choose experimental conditions in which biochemical reactions are well behaved and rate constants are reasonably constant while results are biologically useful.” I suspect a similar situation applies in chemistry and science in general. A scientist does what she/he can do, even if it is much less than what she/he wishes to do. That behavior is expedient, practical, admirable, and a necessary reflection of human limitations, I think.

*Fig.1*

Current Flow is Not the Same in a Sequence of Reactions

 

 

**Sequences of Reactions.** Biochemical reactions occur almost always in sequence in one dimension in a geometry similar to a series combination of elements in an electrical circuit. Fig. 1 is a simple example. Conservation of charge implies that the current is identical in each reaction in the sequence of Fig. 1: the flow of charge is continuous without loss.

Interruptions in current flow in one place in Fig. 1 must be able to interrupt a chemical reaction somewhere else according to Kirchoff’s current law. If current is forced to be zero in one reaction, it must be zero in every reaction, because Kirchoff’s current law says current is the same everywhere in a sequence.

‘Completing the circuit’ is a necessity familiar to all who have wired up a circuit, even if they do not know call the principle ‘Kirchoff’s current law’. The flow of charge is continuous on all scales because conservation of charge and its movement is true on all scales. Interrupting current on an atomic scale (in a semiconductor diode, or an ionic channel) interrupts in macroscopically in the inputs and outputs of the diode, or in the solutions and cells connected to the ionic channel.

Completing the circuit is not necessary in hypothetical sequences of chemical reactions that satisfy **only** conservation of mass because the law of mass action does not guarantee continuity of any flow. Flows can accumulate mass between reactions and satisfy conservation of mass that way, without being continuous. Kirchoff’s current law does not allow charge to accumulate that way when displacement current is included in the definition of current [22, 40, 74], see below.

**Rate constants in the law of mass action are independent of each other.** The rate constants in one reaction of a sequence are independent of other reactions in the chemical kinetic models, for example, in the sequence of Fig. 1. **If rate constants are constant, they cannot depend on each other.** If the masses that flow are charged, and rate constants are independent, the flow of charge cannot be continuous except under very special circumstances. The mathematical proof below shows (see eq. ??) that Kirchoff’s current law is violated except in very special circumstances.

The law of mass action needs to be extended to include charge and charge movement so current flow in a sequence of reactions is equal. After the extension, current will necessarily flow in complete circuits and interrupting current flow in one reaction of a sequence will stop current flow in all of them. With this extension, we can hope to make models of chemical and biochemical reactions as robust as those of our electronic technology.

We turn now to a discussion of charge movement so readers less familiar with electric circuits will know what I am talking about.

**Charge is an abstraction that is conserved.** Physicists teach that conservation of charge is universally true, exact from very small to very large scales and very small times to very large times.**[[1]](#footnote-1)†** What physicists often do not teach clearly is that charge is an abstract idea. Many students believe that charge is always carried by particles and so conservation of mass (of charged particles) implies conservation of charge. That is not true. Displacement current is not carried by the movement of particles. [22, 40, 74] Current flow in semiconductors is carried by quasi-particles. [42] Quasi-electrons and quasi-holes are mathematical fictions. They are not real particles and do not carry current in the same sense that the movement of charged particles carries current. “Hole” this is a quasi-particle, like the electron, but of opposite charge; it corresponds to the absence of an electron from a single particle state that lies just below the Fermi level” consisting of a real electron and an exchange correlation hole …” p. 68 of [41]

We see that charge in one physical system is quite different from charge in another. Charge flow is **not** simply the physical movement of particles of definite charge (and mass). Charge flow in a vacuum capacitor, in an ionic solution, in a wire are all quite different. Current is not just the movement of ions or electrons or protons.

The essential idea of electricity and magnetism as explained in textbooks [40, 74] is that **charge flow is continuous** (without loss) **no matter what the physical nature of the charge.** Fig. 2 tries to show this idea in a concrete way that anyone can build in their lab. It shows time‑varying currents in a sequence of different devices each with different physics and different constitutive laws.

**Ionic Conductor.** An ionic conductor is shown in Fig. 2 as a cylinder containing Na+Cl−. Here, current flow (at a frequency say of 1 Hz) is almost entirely the physical movement of charged particles, of ions, say sodium and chloride ions in cylinder and follows ideal constitutive laws that are quite simple when concentrations are low but flows are not large enough to change concentrations. These ions are nearly hard spheres, as detailed in texts of physical chemistry and that finite size makes constitutive equations much more difficult. The finite size of these spheres is significant at concentrations greater than say 2×10−2 molar in the ionic mixtures found everywhere in biology.

Numerical difficulties in dealing with spheres are substantial. Spheres must be computed in three dimensions because spheres do not exist in one and two dimensions. Computation of the forces that prevent overlap of spheres is difficult because those forces are strong and vary steeply. They are singular in the sense of numerical mathematics.

Fig. 2

*redrawn from reference [16]*

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| Charge Is Abstract, with Different Physics In Different Systems |
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**Vacuum Capacitor.** Let’s move to a vacuum capacitor, in which the space between the two plates is completely empty of matter (as it would be in outer space, for example). The current flow through this capacitor is just as real as the movement of ions of Na+Cl− in the cylinder even though no particles or spheres are present, and no mass moves at all.

The displacement current between the plates of the capacitor is a property of the electric field itself, as explained in textbooks of electricity and magnetism (e.g., [22, 40, 74]) and is described by an exact and simple constitutive equation  where the displacement current  (amps) is strictly proportional to the capacitance  (farads) and the time rate of change of the voltage across the capacitor. This displacement current induces a magnetic field just as current carried by ions produces a magnetic field. Indeed, without displacement current in a vacuum, Maxwell’s equations cannot account for the propagation of sunlight through the vacuum of space. With displacement current, light propagation is a solution of the equations, and in fact the speed of propagation of light can be computed from measurements of electrical and magnetic constants, entirely independent of measurements on light itself.

**Dielectric Capacitor.** In the dielectric capacitor in Fig. 2 (filled with real material, for example the plastic = Polytetrafluoroethylene = PTFE = Teflon), current flow is more complex, and involves the effect of an applied electric field on the spatial distribution of the electric charge intrinsic to the atoms, molecules, and substance of the dielectric. (‘Intrinsic’ here means the distribution of charge present when there is no applied electric field. Note that the electric charge varies significantly on all scales from subatomic to macroscopic in different parts of Fig. 2)

The properties of dielectrics cannot be described in detail here because they vary so much with material, voltage, and time. But some properties need emphasis. Intrinsic charges do **not** move from plate to plate in a real capacitor. The current that flows from plate to plate and within the dielectric is a displacement current not carried by the movement of mass (any substantial distance). The current in a real capacitor is an abstraction. It is the sum of the vacuum displacement current and the material displacement current (in the dielectric) produced by the distortion of intrinsic charges. Calculating this material displacement current involves solving a nonequilibrium time dependent version of the Schrödinger equation, a difficult task.

Scientists have avoided the difficulties of solving the Schrödinger equation by using approximations. They discuss an ideal dielectric with properties independent of field strength and independent of time. The approximation over a wide range of electric field strength is satisfactory in most materials. The approximation over a wide range of times is poor. For example, most solutions of ions in water have effective dielectric coefficients that change—after a step electric field is applied—from about 2 to about 80 as time moves from 10‑15 sec to say 10‑3 sec. A factor of 40 is not a small change, but biology is controlled by atomic structures that move significantly in 10‑15 sec and have dramatic effects on biological functions that take 10‑3 sec or much longer!

Electric fields applied to dielectrics are often strong in experiments and in the course of biological function, during a biological action potential, some 0.1 volts across a 2×10-9 meter thick membrane or a 3×10-10 meter selectivity filter (EEEE group in a calcium channel), 5×107 to 3×108 volts/meter. These electric fields distort the distribution of charge within the dielectric on many scales of time and distance, including electrons inside molecules and atoms, They reorient polar molecules that have an intrinsic asymmetric distribution of charge. Charges only move a small amount in dielectrics in response to the applied electric field—–reminiscent of the sloshing of tides in the ocean and on beaches on the earth created by the moon’s gravitational field— **and they return to their resting position when the field is turned off** but those small movements of charge produce large effects because the electric field is so strong.

**Any Electric field is extraordinarily strong.**  **A one per cent deviation in the density of (net) charge has enormous effects.** As Feynman mentions unforgettably [22] in his introductory paragraphs about the electric field, **one per cent of the charge in a person at arm’s length produces force enough to lift the earth!** Note that the electric charge in different devices of Fig. 2 varies on very different scales from subatomic (in dielectrics) to atomic (in diodes, ionic conductors, etc.) to macroscopic and so treatments of charge must be multiscale.

**Concentration effects are very small.** Diffusion forces are tiny perturbations on the electric force field and energies involved in diffusion are tiny perturbations in the total energy of charged systems. A one percent deviation in density of mass has tiny effect on diffusion and chemical reactions, hardly noticeable. Electric forces and energies are not small parts of the total energy and it is not wise—to say the least—to treat them as perturbations of uncharged systems.

Charge is indeed an abstraction. In each of the devices in Fig. 2, charge follows different laws, because it has different physical properties, sometimes carried by charged particles, sometimes produced by the rate of change of electric field (displacement current), sometimes the ‘movement’ of quasi-particles, sometimes the ‘movement’ of electrons in a macroscopically delocalized quantum state of a wire.

Charge always flows without loss in each device. That abstract property of charge movement is always true, but the ‘laws’ describing current as a function(al) of time and potential (for example) depend on the physical nature of the charge and its movement. Different devices also have different relationships between current, voltage, and time. These different ‘constitutive’ equations described in the engineering literature in great detail.

The different constitutive equations combine with Kirchoff’s current law to describe current flow from one place and one device to another. Together, the equations describe the universal fact that interrupting current flow in one reaction (of a sequence of reactions) will interrupt current flow in every other reaction (in that sequence) even if one reaction occurs on the atomic scale (scale = 1 Å) and is meters away from the interruption. One is reminded of the effects of unplugging a computer. Interruption of a circuit meters away from the diodes of the computer’s power supply stops the flow of holes and semi-electrons across atomic scale junctions of semiconductor diodes, often of the ***PN*** variety. The electric field changes so strongly in response to the interruption (because of the circuit properties of the input of the computer power supply) that the electric field increases dramatically, and exceeds the dielectric strength of air. Electrons are stripped off the atoms of the air, a plasma is created, and we can see the light produced as a spark, and perhaps be frightened by the power of the electric field. After all, the spark could start a fire, in nearby upholstery or drapes.

The electric field of charge must be respected on all scales because it has such large effects on all scales. Although charge is an abstraction in each device, it is conserved ‘perfectly’ as charge moves from device to device, according to Kirchoff’s current law. These properties are abstract but they are also concrete. The continuity of current flow is exploited continually in the design of integrated circuits [40] and the careful handling of current in grounds is essential for the smartphone, imaging, and global positioning devices we use every day. If Kirchoff’s current law were ignored, the devices would not work.

**Vacuum tubes.** In the next device in Fig. 2, we consider charge movement in a vacuum tube diode. Vacuum tubes control the flow of electron current by changes in their internal electric fields and were called ‘valves’ in the UK for that reason. Ion channels can control the flow of ionic current the same way. Both follow simple constitutive laws of rectifiers, as described in textbooks of electronic devices or ion channels.

In vacuum tubes, current is indeed carried by a stream of isolated charged particles, electrons with a definite mass and charge, moving through a vacuum, interacting only through their electric fields. At 1 Hz, essentially all the current in a vacuum diode is carried this way. Current through a diode is not proportional to the voltage across the diode because the electric fields within the diode change shape, despite the simple physics of conduction. The fields change shape as the voltage across the diode is changed because of screening and shielding. The electric fields within the tube are different at small and large potentials. The different internal electrical fields change the current flow, creating rectification. The electric field creates a large barrier in one direction so current in that direction is small; the electric field creates a small barrier in the other direction, and current in that other direction is large. Rectification is of considerable historical interest, because vacuum tubes allowed the early detection of radio waves in American’s homes in the 1930’s. The demand for portable radios led to solid state diodes, then solid state ‘triodes’, field effect transistors, integrated circuits, and our modern digital world.[64]

Vacuum tube diodes had certain difficulties and were quickly replaced with semiconductor diodes, the next device in Fig. 2. Here current is carried by quasi-particles [42], called holes and ‘electrons’ (better named quasi-electrons, or semi-electrons in my view) that simplify the [quantum mechanical](http://en.wikipedia.org/wiki/Quantum_mechanics) [many-body problem into the](http://en.wikipedia.org/wiki/Many-body_problem) much simpler motion of imagined quasiparticles. These quasi-electrons and holes are defined because they interact much more simply than the totality of real electrons in the semiconductor.[42, 58, 60, 69, 70] To quote a recent textbook “Electron is a quasi-particle consisting of a real electron and an exchange correlation hole … a cloud of effective charge of opposite sign due to exchange and correlation effects arising from interactions with **all** other electrons of the system” p. 68 of [41], my emphasis. The quasi-particles are described by one electron mean field models, both simulations and theories, for example, the Poisson drift diffusion equations often called *PNP* (for Poisson Nernst Planck) in biophysics and nanotechnology. The quasi-particles are quite distinct from real particles and move differently.

**Charge carriers in semiconductors can be really quite strange** from a physical point of view. Charge carriers in silicon and germanium do not exist outside the lattice of the semiconductor as distinct entities. They do not exist in the same sense that Na+ and Cl− ions exist, but are mathematical representations, with lifetimes sometimes as short as milliseconds. Quasi-particles are the second derivative of a Fermi surface of silicon and germanium semiconductors, under particular conditions. They are defined, as mentioned above, to allow easy analysis compared to the intractable quantum mechanical many-body problem of a macroscopic semiconductor with atomic scale diode (e.g., *PN* ) junctions driven by far field boundary conditions.

Much of the success of our semiconductor, digital, and video technology is due to the accuracy of the constitutive *PNP* equations describing holes and semi-electrons. They describe the characteristics of semiconductor devices of many different types, as different as an exponentiator and a logarithmic convertor.

*PNP* is so useful because almost all the devices of our digital technology work in a restricted set of conditions in which flows are crucial but are of a special type. Almost all of these devices use power supplies to maintain different voltages at different locations far from the *PN* junctions of the device itself (and thus require a global treatment of the electric field). These voltages perturb the distribution of velocities of charged particles so the distribution has net flow [14, 20]. The slight perturbation is enough to imply the *PNP* equations and those provide enough nonlinearity to make amplifiers, switches, and the full set of logical circuits necessary to make a computer.[58, 60, 69, 70]

***PNP* equations describe a wide variety of current voltage relations and devices.** The *PNP* equations are the constitutive laws used to describe semiconductor devices as their current voltage relations change drastically (with voltage, for example) from that of a linear amplifier, to a switch, to an exponentiator, a multiplier or even a logarithmic converter. [58, 60, 69, 70] Nonlinear input output relations as varied as these enable a rich variety of devices. Nonlinear input output relations as diverse as these are not described easily—or described at all, for that matter—in most areas of physics and chemistry.

The descriptive power of *PNP* is so important that the intrinsic physical properties of semiconductor devices (e.g., concentration of fixed charge dopants, geometry, and recently even dielectric coefficients) are scaled by their designers so *PNP* remains a good description, even as the size of the device is decreased in successive iterations of Moore’s law,[52, 55, 56] by factors of a billion or so over 50 years. Evidently, it is more important for the designers (and marketplace) that an equation describes behavior robustly over a range of conditions than that the device be as fast or small as possible. [8-10, 66] Evolutionary selection in biology also seems to choose robustness over efficiency in many cases. Devices of nanotechnology need to be robust, I believe, before they will be used extensively.

***PNP* is not enough, however, when ionic solutions are involved.** A great deal of effort has been spent applying *PNP* equations to electrochemical systems [4-7, 12, 18, 19, 32, 72] hoping they might serve as adequate constitutive equations, but that is not the case. The nonideality of ionic solutions, arising in large measure from the finite size of ions, demands more powerful mathematics than the partial differential equations of *PNP* used in computational electronics.

**Our next device is a resistor**, which in some ways is the easiest to describe, because current is proportional to voltage with a single proportionality constant over a wide range of voltages, times, and conditions. A resistor follows Ohm’s constitutive law with a resistance independent of potential over a wide range from 10-5 volts to say 100 volts, and fromvalues of resistance from 10-1 ohm to 108 or 109 ohms.

The range of validity of Ohm’s law is an enormous help in circuit design. Circuit models involving resistors, capacitors, inductors and operational amplifiers are transferrable. They behave as real devices behave without changes in parameters. Largely for that reason, designs are inexpensive and robust.

Despite the simplicity of resistors, the actual current carrier in a carbon resistor, is unclear, at least to me. No one cares very much I suspect, because the device works nearly perfectly. The carrier of charge does not matter very much. What matters is the constitutive law that describes the relation of current, voltage and time. The constitutive law satisfies conservation of charge. It can combine conservation of charge and matter with the properties of the physical constituent if written with the energetic form of the variational calculus.

**Current flow in wires.** The current carriers in a wire are delocalized electrons in the simple case of a single solid conductor of metal, and follow the simplest constitutive law of all at long(ish) times, say times longer than 10-5 sec.

But most of our electronics occurs at times much shorter than that. Here we confront the importance of the time variable. **The physical nature of current flow depends on the time scale on which the current changes even in wires.** The range of time scales in our technology is enormous from more than one second to less than 10−9 sec. At those times, the wires that must be used are are often ***TWISTED*** pairs, each made itself of many very fine wires. Twisting is. Without twisting. these pairs of wires do not allow successful connections to the internet because without twisting. rapidly changing signals are not carried reliably by wires. At those short times, currents flow outside wires, guided by the conductor, to be sure, but outside the conductor nonetheless [40]. The twisting of wires is a necessity if they are to carry signals robustly and reliably so we can use them in our video devices and smartphones.

**Physical nature of current depends on time scale.** The physical nature of current in almost any system depends on the time scale, and differs at different times as much as it differs in different devices. Constitutive equations depend on time. Different devices have different constitutive equations with different time dependence.

Again, charge movement is an abstraction, different at different times in one device, as well as different in different devices.

**Charge Movement in Batteries.** Batteries are present in Fig. 2 both as an isolated device and at the Ag||Ag+Cl− interface between wires and Na+Cl− solution in the conducting cylinder previously discussed. I hesitate to describe a constitutive equation for the flow of charge in these systems because it is so very complex, different in different devices, and its interaction with surface charges is also subtle and important. It is enough to say here that current flow through electrochemical systems is carried by a wide range of charge carriers. Constitutive equations of different electrochemical systems differ and change nature dramatically with time and frequency. At short times, at 10-6 sec—that are still long and slow compared to the times important in computers—current from the AgCl wire into the Na+Cl− solution is entirely displacement current lagging behind voltage but at the longer times characteristic of biological systems (greater than say 0.1 sec) the current is carried by a complex combination of Ag+ and Cl− ions with negligible displacement current.

**We conclude that charge is indeed an abstraction** with different physical meanings in different systems and at different times. No one can visualize and no one knows—at least I do not know anyone who knows—why or even how this abstraction can be so perfectly conserved under all conditions and on all scales, from Angstroms to meters, from femtoseconds to seconds. The reader may have difficulty visualizing the interactions that enforce conservation of the abstraction ‘charge’ in all these devices with all these properties over the entire time scale. I certainly do. However,

**Charge is conserved on and between all scales in all conditions,** even if we cannot visualize how. Experiments demonstrate that fact. Current does flow continuously without loss in a circuit. Interruptions in current flow meters from a chemical reaction does interrupt current flow in that reaction, as can easily be demonstrated by disconnecting a car battery, for example (be careful!). Abstract charge is conserved ‘exactly’ even if we cannot visualize how that happens.

Science often contains mysteries that cannot be understood and (e.g., why is the charge on an electron 1.6 ×10−19 coulombs?) and explanations that cannot be visualized. Consider Maxwell’s attempts to visualize his equations as properties of an ether. We wonder and move on, hoping our successors can do better than we have.

Biologists and engineers cannot function if they linger on the mysteries they do not understand. Think of Thomas Henry Huxley looking at the shortening of the striations of muscle [35] that were not understood until his grandson studied them many years later.[34] Think of Lee DeForest using vacuum tubes without understanding how they work. Biologists and engineers cannot afford to wait to understand everything. They must isolate the mysteries and move on to study other things. Here, we move on to discuss devices and the theories and simulations used to understand them.

**Devices, effective parametes, and transferrable theories.** Theories and simulations that do not conserve charge are not likely to be generally useful. At best they will be incomplete.

Parameters of devices can often be chosen so an incomplete theory or simulation describes a system in one set of conditions. Chemistry and biology are filled with examples. Chemical reactions follow rate equations, but the rates are not constant, as conditions change, even though theory assumes they should be. Biology describes enzymes with one set of parameters but finds those are changed when conditions change and attributes that somewhat mysteriously to ‘allosteric effects’ and conformation changes.

Biology and much of chemistry works under a wide range of conditions and so incomplete theories with effective parameters have limited use. Even if sensible, even if valid, theories (and simulations) with effective parameters like these are not accurate enough to design of robust devices. By leaving out something important, those theories or simulations leave out an energy term that is almost certain to vary with conditions. The resulting effective parameters change in large and unpredictable ways.

Theories and parameters of incomplete theories cannot be transferred successfully from one set of experiments and conditions to another. Incomplete theories are not likely to be transferable in the language of the chemistry literature. Devices designed from incomplete theories or simulations are unlikely to be robust or work well under a range of conditions.

Biological systems analyzed with non-transferable theories (or simulations) are unlikely to be realistic in general because biological systems almost always work in a range of ionic concentrations different from those used in the laboratory. Biological systems usually work in mixtures with a range of Ca2+ concentrations, in which Ca2+ concentration has important practical effects, often turning systems on or off or controlling their rate monotonically. Simulations and theories in biology have limited use until they are calibrated so we can be sure they actually are correct in the range of conditions and Ca2+ concentrations the biological system uses. Simulating Ca2+ activity in pure solutions is a challenging problem. [67] Simulating Ca2+ activity in biological mixtures in the 10-8 to 2 × 101 molar concentrations that are physiological, has not been attempted as far as I know.

Biological and chemical science will benefit enormously if theories and simulations can be made transferable, using one set of parameters to describe systems in a range of conditions, as many physical and most engineering theories and simulations do. I believe the law of mass action must be extended to conserve charge before theories and simulations can be made transferable from condition to condition, from physics, to chemistry to biology, using only mathematics.

**How to extend the law of mass action?** An obvious way to extend the law of mass action is to include an electrical potential in the rate constant to ‘right the rates’ by making them exponential functions of potential. This in fact has been done for a very long time in the study of reactions at the electrodes of electrochemical cells. The Butler-Volmer and Tafel equations [2] include electrical potential in rate constants in an empirical way with limited [27, 28] but real success.

 The difficulties of just embedding an electrical potential in a rate constant are formidable. **We must ‘fix the fields’ so they are global and depend on current flow everywhere.** The only way to ‘right the rates’ is to fix the fields, **everywhere.** Otherwise, they cannot conserve charge and support continuity of current as required by Kirchoff’s current law.

1. A thermodynamic treatment is clearly impossible since the goal is to calculate large fluxes and currents that do not occur in a thermodynamic system at equilibrium by definition without flows. Clearly a frictional treatment of flux over a large potential barrier as in the classical drunken man problem of Kramers [23, 29] is a necessary step forward, even if it is inconsistent because it does not compute the potential barrier from the charges in the system. A simple generalization of Kramers for any shape barrier is available [14] and needed [1] because so few barriers are both symmetrical and large as required in classical high barrier approximations.
2. A general theory of all nonequilibrium processes is not likely to be useful: a general theory has to describe too much. A general theory must include hydrodynamic behavior of considerable complexity, since aqueous solutions are fluids satisfying the Navier Stokes equations of fluid mechanics. A general theory should also include explosions since they occur with regrettable frequency at electrodes of electrochemical cells, when H2 gas is generated (inadvertently) by an overvoltage.
3. The rate constants in one reaction must depend on the electrical potential in far locations. Otherwise, interrupting current flow in a far location is not likely to interrupt current locally. This requirement implies that the electrical potential must be determined by a global equation like Poisson’s equation, including boundary conditions far from the individual chemical reactions.
4. The rate constants in one reaction must depend on concentrations throughout the system because the solutions containing the reactants are not ideal. A general characteristic of nonideal solutions is that ‘everything depends on everything else’. More specifically, the activity of one reactant (the free energy per mole) depends on the concentrations of other species in practice, as well as in principle.

**Variational methods can extend the law of mass action.** If the goal is to build transferable systems, so we can build robust devices, as in electronic technology, we must use a mathematics that allows interactions of charges and fields, extending from atomic to macroscopic scales.

Variational methods are designed to deal with systems with multiple forces and flows, in which interactions are unavoidable and complex. In these systems interactions must be included in all analysis. Otherwise, theories have more adjustable parameters than can be determined experimentally and still cannot deal with a range of conditions, because interactions change with conditions in ways too complex for ordinary theories.

Sadly, theories of ionic solutions—even the assignment of properties to single ions in infinitely dilute solutions [33]—seem to have these difficulties. Theories of ionic solutions need a large number of adjustable parameters and still cannot describe biological solutions (for example). References [13, 15, 25, 26, 43-46, 59, 68, 75] draw particular attention to the difficulties.

The more powerful tools of variational calculus are needed, in my view, so the mathematics used automatically deals with interactions that vary dramatically with conditions. If the mathematics does not deal with interactions, those interactions will not be computed correctly and will wreck theories, based (for example) on algebraic descriptions of interactions or theories based on the law of mass action with constant rate constants.

Energetic variational methods are particularly useful because they allow multiscale derivation of partial differential equations (and far field boundary conditions) from physical principles when multiple fields are involved, like convection, diffusion, and migration in an electric field. Energetic variational principles have recently become available for systems involving friction [24, 36, 73], that is to say, for systems involving ionic solutions [17, 32, 57, 71]. They combine the full power of the Navier Stokes equations (1) with either a Lennard Jones representation of finite size ions or (2) with a density functional theory of ionic solutions built from Rosenfeld’s DFT of liquids [17, 37]. These formulations must be computed in three dimensions because spheres do not exist in one and two dimensions. These theories and their simplifications [31, 47, 48] are difficult to compute in three dimensions because of the steeply singular forces used to ensure that atoms do not overlap. Overlap must not be permitted because spheres cannot overfill space: space can be saturated with spheres. Saturation effects are a main cause of nonideality, particularly in the extremely crowded conditions in and near enzyme active sites, ionic channels, nucleic acids, and the working electrodes of electrochemical cells where twenty molar solutions are not uncommon [39]. (As a rule of thumb, ions are crowded where they are most important in technology and biology.)

If saturation is described by a Fermi-like distribution—as recently derived for spheres of unequal sizes [49, 50]—some of these difficulties can be attacked. A fourth order partial differential equation can be derived [51] which is easily integrated in three dimensions, after it is reduced to a pair of second order partial differential equations (with carefully defined boundary conditions) and computed with appropriate numerical methods.

But it is still not clear how best to apply any of these methods to chemical reactions (Fig. 1) described by the law of mass action with rate constants extended to be functions or functionals and not constants.

**And that is not a bad place to move this essay towards its coda.** We have the tools, and we know the goal—a global treatment combining conservation of mass locally in chemical reactions with conservation of charge locally and globally. Now we must do the work.

**Coda: our grasp must be sure, but our reach should exceed our grasp**, as we do the work needed to understand more. We must grasp both charge conservation and mass action before we can produce robust theories (or simulations) of chemical reactions in ionic solutions that successfully use one set of parameters in a range of conditions, and include the global properties of the electric field.

There is no engineering without numbers and computations. All calculations from theories and simulations—of electronic or ionic or biologically inspired devices—must be checked and calibrated against known results. Otherwise devices built from those calculations will not work.[11, 53, 61] There is little engineering without robust theories because devices have much less tolerance for error than abstractions. If theories and simulations of electrical devices were not robust, if parameters had to be changed as conditions changed, our electronic technology would be severely limited, to say the least.

Chemical reactions involving charge must be within our theoretical grasp before we can develop transferable theories, in my view. Only then can we build devices that perform as expected, as electronic devices usually do. Only then can we expect exponential growth in molecular engineering. whether biological or technological. It seems no coincidence that exponential growth in electronic technology came after scientists had a secure grasp of global electrostatics and the *PNP* equations of electrodiffusion in semiconductors. Let us hope that energetic variational methods can grasp ions and chemical reactions as well and as successfully as *PNP* has grasped the useful properties of holes and (semi)-electrons in silicon and germanium.

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**References**

1. Barcilon, V., D. Chen, R.S. Eisenberg, and M. Ratner, Barrier crossing with concentration boundary conditions in biological channels and chemical reactions*.* J. Chem. Phys., 1993. **98**: p. 1193–1211.

2. Bard, A.J. and L.R. Faulkner, *Electrochemical Methods: Fundamentals and Applications.* . 2nd Edition ed. 2000, New York: John Wiley & Sons.

3. Barratt, J.-L. and J.-P. Hansen, *Basic concepts for simple and complex liquids*. Cambridge University Press, Mar 1, 2003 - Science - 296 pages. 2003: Cambridge University Press. 296.

4. Bazant, M.Z., K. Thornton, and A. Ajdari, Diffuse-charge dynamics in electrochemical systems*.* Physical Review E, 2004. **70**: p. 021506.

5. Boda, D., Monte Carlo Simulation of Electrolyte Solutions in biology: in and out of equilibrium*.* Annual Review of Compuational Chemistry, 2014. **10**: p. 127-164.

6. Burger, M., Inverse problems in ion channel modelling*.* Inverse Problems, 2011. **27**(8): p. 083001.

7. Coalson, R.D. and M.G. Kurnikova, Poisson-Nernst-Planck theory approach to the calculation of current through biological ion channels*.* IEEE Trans Nanobioscience, 2005. **4**(1): p. 81-93.

8. Critchlow, D.L., MOSFET Scaling-The Driver of VLSI Technology*.* Proceedings of the IEEE, 1999. **87**(4): p. 659-667.

9. De Loach, B.C., Jr., The IMPATT story*.* Electron Devices, IEEE Transactions on, 1976. **23**(7): p. 657-660.

10. Dennard, R.H., F.H. Gaensslen, V.L. Rideout, E. Bassous, and A.R. LeBlanc, Design of ion-implanted MOSFET's with very small physical dimensions*.* Solid-State Circuits, IEEE Journal of, 1974. **9**(5): p. 256-268.

11. Eisenberg, B., Multiple Scales in the Simulation of Ion Channels and Proteins*.* The Journal of Physical Chemistry C, 2010. **114**(48): p. 20719-20733.

12. Eisenberg, B., *Crowded Charges in Ion Channels*, in *Advances in Chemical Physics*, S.A. Rice, Editor. 2011, John Wiley & Sons, Inc.: New York. p. 77-223 also on the arXiv at <http://arxiv.org/abs/1009.1786v1001>.

13. Eisenberg, B., Life’s Solutions are Not Ideal*.* Posted on arXiv.org with Paper ID arXiv:1105.0184v1, 2011.

14. Eisenberg, B., Mass Action in Ionic Solutions*.* Chemical Physics Letters, 2011. **511**: p. 1-6.

15. Eisenberg, B., Ionic Interactions Are Everywhere*.* Physiology, 2013. **28**(1): p. 28-38.

16. Eisenberg, B., Shouldn’t we make biochemistry an exact science? ASBMB Today, 2014. **13**(9:October): p. 36-38.

17. Eisenberg, B., Y. Hyon, and C. Liu, Energy Variational Analysis EnVarA of Ions in Water and Channels: Field Theory for Primitive Models of Complex Ionic Fluids*.* Journal of Chemical Physics, 2010. **133**: p. 104104

18. Eisenberg, R. and D. Chen, Poisson-Nernst-Planck *(PNP)* theory of an open ionic channel*.* Biophysical Journal, 1993. **64**: p. A22.

19. Eisenberg, R.S., *Atomic Biology, Electrostatics and Ionic Channels.*, in *New Developments and Theoretical Studies of Proteins*, R. Elber, Editor. 1996, World Scientific: Philadelphia. p. 269-357. Published in the Physics ArXiv as arXiv:0807.0715.

20. Eisenberg, R.S., M.M. Kłosek, and Z. Schuss, Diffusion as a chemical reaction: Stochastic trajectories between fixed concentrations. J. Chem. Phys., 1995. **102**: p. 1767-1780.

21. Evans, L.C., *An Introduction to Stochastic Differential Equations*. 2013: American Mathematical Society. 150.

22. Feynman, R.P., R.B. Leighton, and M. Sands, *The Feynman: Lectures on Physics, Mainly Electromagnetism and Matter*. Vol. 2. 1963, New York: Addison-Wesley Publishing Co., also at <http://www.feynmanlectures.caltech.edu/II_toc.html>. 592.

23. Fleming, G. and P. Hänggi, *Activated Barrier Crossing: applications in physics, chemistry and biology*. 1993, River Edge, New Jersey: World Scientific.

24. Forster, J., *Mathematical Modeling of Complex Fluids*, in *Department of Mathematics*. 2013, University of Wurzburg: Wurzburg, Germany. p. 67.

25. Fraenkel, D., Simplified electrostatic model for the thermodynamic excess potentials of binary strong electrolyte solutions with size-dissimilar ions*.* Molecular Physics, 2010. **108**(11): p. 1435 - 1466.

26. Fraenkel, D., Computing Excess Functions of Ionic Solutions: The Smaller-Ion Shell Model versus the Primitive Model. 1. Activity Coefficients*.* Journal of Chemical Theory and Computation, 2014.

27. Gutman, E.M., Can the Tafel equation be derived from first principles? Corrosion Science, 2005. **47**(12): p. 3086-3096; note IMPORTANT correction doi:3010.1016/j.corsci.2006.3008.3001 Corrosion Science 3048 (2006) 3886.

28. Gutman, E.M., Corrigendum to “Can the Tafel equation be derived from first principles?” [Corrosion Science 47 (2005) 3086–3096]*.* Corrosion Science, 2006. **48**(11): p. 3886 IMPORTANT correction to <http://dx.doi.org/3810.1016/j.corsci.2005.3805.3055>

29. Hänggi, P., P. Talkner, and M. Borokovec, Reaction-rate theory: fifty years after Kramers. Reviews of Modern Physics, 1990. **62**: p. 251-341.

30. Hansen, J.-P. and I.R. McDonald, *Theory of Simple Liquids*. Third Edition ed. 2006, New York: Academic Press. 428.

31. Horng, T.-L., T.-C. Lin, C. Liu, and B. Eisenberg, PNP Equations with Steric Effects: A Model of Ion Flow through Channels*.* The Journal of Physical Chemistry B, 2012. **116**(37): p. 11422-11441.

32. Hsieh, C.-y., Y. Hyon, H. Lee, T.-C. Lin, and C. Liu, Transport of charged particles: entropy production and maximum dissipation principle*.* Available on <http://arxiv.org/> as 1407.8245v1.

33. Hünenberger, P. and M. Reif, *Single-Ion Solvation. Experimental and Theoretical Approaches to Elusive Thermodynamic Quantities.* 2011, London: Royal Society of Chemistry. 690.

34. Huxley, A.F., *Looking Back on Muscle*, in *The Pursuit of Nature*, A.L. Hodgkin, Editor. 1977, Cambridge University Press: New York. p. 23-64.

35. Huxley, T.H., *Crayfish, an Introduction to the Study of Biology*. 1884, London: K. Paul, Trench. 371.

36. Hyon, Y., D.Y. Kwak, and C. Liu, Energetic Variational Approach in Complex Fluids : Maximum Dissipation Principle*.* Discrete and Continuous Dynamical Systems (DCDS-A), 2010. **26**(4: April): p. 1291 - 1304.

37. Hyon, Y., J.E. Fonseca, B. Eisenberg, and C. Liu, Energy variational approach to study charge inversion (layering) near charged walls*.* Discrete and Continuous Dynamical Systems Series B (DCDS-B), 2012. **17**(8): p. 2725 - 2743.

38. Jaffe, R.L., The Casimir Effect and the Quantum Vacuum*.* Available on <http://arxiv.org/> as <http://arxiv.org/abs/hep-th/0503158>, 2005.

39. Jimenez-Morales, D., J. Liang, and B. Eisenberg, Ionizable side chains at catalytic active sites of enzymes*.* European Biophysics Journal, 2012. **41**(5): p. 449-460.

40. Joffe, E.B. and K.-S. Lock, *Grounds for Grounding*. 2010, NY: Wiley-IEEE Press. 1088.

41. Kaxiras, E., *Atomic and Electronic Structure of Solids*. 2003, New York: Cambridge University Press. 676.

42. Kittel, C., *Solid-State Physics, Eighth Edition*. 8th Edition ed. 2004, New York: Wiley. 704.

43. Kontogeorgis, G.M. and G.K. Folas, *Thermodynamic Models for Industrial Applications: From Classical and Advanced Mixing Rules to Association Theories*. 2009: John Wiley & Sons. 721.

44. Kraus, C.A., The present status of the theory of electrolytes*.* Bull. Amer. Math. Soc., 1938. **44**: p. 361-383.

45. Kunz, W., *Specific Ion Effects*. 2009, Singapore: World Scientific 348

46. Laidler, K.J., J.H. Meiser, and B.C. Sanctuary, *Physical Chemistry*. Fourth ed. 2003: BrooksCole, Belmont CA. 1060.

47. Lee, C.-C., H. Lee, Y. Hyon, T.-C. Lin, and C. Liu, New Poisson–Boltzmann type equations: one-dimensional solutions*.* Nonlinearity, 2011. **24**(2): p. 431.

48. Lin, T.C. and B. Eisenberg, A new approach to the Lennard-Jones potential and a new model: PNP-steric equations*.* Communications in Mathematical Sciences, 2014. **12**(1): p. 149-173.

49. Liu, J.-L., Numerical methods for the Poisson–Fermi equation in electrolytes*.* Journal of Computational Physics, 2013. **247**(0): p. 88-99.

50. Liu, J.-L. and B. Eisenberg, Correlated Ions in a Calcium Channel Model: A Poisson–Fermi Theory*.* The Journal of Physical Chemistry B, 2013. **117**(40): p. 12051-12058.

51. Liu, J.-L. and B. Eisenberg, Poisson-Nernst-Planck-Fermi theory for modeling biological ion channels*.* J Chem Phys, 2014. **141**(22): p. -.

52. Lundstrom, M., Applied Physics Enhanced: Moore's Law Forever? Science, 2003. **299**(5604): p. 210-211.

53. Maginn, E.J., From discovery to data: What must happen for molecular simulation to become a mainstream chemical engineering tool*.* AIChE Journal, 2009. **55**(6): p. 1304-1310.

54. McQuarrie, D.A., *Statistical Mechanics*. 1976, New York.: Harper and Row.

55. Moore, G.E., Cramming more components onto integrated circuits*.* Electronics Magazine., 1965. **38**: p. 114–117.

56. Moore, G.E. *Lithography and the future of Moore's law*. 1995. Santa Clara, CA, USA: SPIE.

57. Mori, Y., C. Liu, and R.S. Eisenberg, A model of electrodiffusion and osmotic water flow and its energetic structure*.* Physica D: Nonlinear Phenomena, 2011. **240**(22): p. 1835-1852.

58. Pierret, R.F., *Semiconductor Device Fundamentals*. 1996, New York: Addison Wesley.

59. Pitzer, K.S., *Thermodynamics*. 3rd ed. 1995, New York: McGraw Hill. 626.

60. Plummer, J.D., M.D. Deal, and P.B. Griffin, *Silicon VLSI Technology: Fundamentals, Practice, and Modeling* 2000, New York: Prentice Hall. 817 pages

61. Post, D.E. and L.G. Votta, Computational Science Demands a New Paradigm*.* Physics Today, 2005. **58**: p. 35-41.

62. Reynaud, S. and A. Lambrecht, Casimir forces*.* available on <http://arxiv.org/> as <http://arxiv.org/abs/1410.2746>, 2014.

63. Rice, S.A. and P. Gray, *Statistical Mechanics of Simple Fluids*. 1965, New York: Interscience (Wiley). 582.

64. Riordan, M. and L. Hoddeson, *Crystal Fire*. 1997, New York: Norton.

65. Rowlinson, J.S., *The Perfect Gas*. 1963, New York: Macmillan. 136.

66. Sah, C.-T., Evolution of the MOS transistor-from conception to VLSI*.* Proceedings of the IEEE, 1988. **76**(10): p. 1280-1326.

67. Saxena, A. and A.E. García, Multisite Ion Model in Concentrated Solutions of Divalent Cations (MgCl2 and CaCl2): Osmotic Pressure Calculations*.* The Journal of Physical Chemistry B, 2014. **119**(1): p. 219-227.

68. Sengers, J.V., R.F. Kayser, C.J. Peters, and H.J. White, Jr., *Equations of State for Fluids and Fluid Mixtures (Experimental Thermodynamics)* 2000, New York: Elsevier. 928.

69. Shur, M., *Physics of Semiconductor Devices*. 1990, New York: Prentice Hall. 680.

70. Vasileska, D., S.M. Goodnick, and G. Klimeck, *Computational Electronics: Semiclassical and Quantum Device Modeling and Simulation*. 2010, New York: CRC Press. 764.

71. Wu, H., T.-C. Lin, and C. Liu, On transport of ionic solutions: from kinetic laws to continuum descriptions*.* available on <http://arxiv.org/> as 1306.3053v2, 2014.

72. Wu, H., T.-C. Lin, and C. Liu, Diffusion Limit of Kinetic Equations for Multiple Species Charged Particles*.* Archive for Rational Mechanics and Analysis, 2015. **215**(2): p. 419-441.

73. Xu, S., P. Sheng, and C. Liu, An energetic variational approach to ion transport*.* Communications in Mathematical Sciences, 2014. **12**(4): p. 779–789 Available on arXiv as <http://arxiv.org/abs/1408.4114>.

74. Zangwill, A., *Modern Electrodynamics*. 2013, New York: Cambridge University Press. 977.

75. Zemaitis, J.F., Jr., D.M. Clark, M. Rafal, and N.C. Scrivner, *Handbook of Aqueous Electrolyte Thermodynamics*. 1986, New York: Design Institute for Physical Property Data, American Institute of Chemical Engineers.

1. **†** Even the Casmir effect of quantum physics [62] is seen as a property of Maxwell’s equations by some [38]. [↑](#footnote-ref-1)