

Ionic Channels in Biological Membranes: Natural Nanotubes Described by the Drift-Diffusion Equations

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An important class of biological molecules—proteins called ionic channels—conduct ions (like Na^+ , K^+ , Cl^-) through a narrow tunnel of fixed charge ('doping'). Ionic channels are the main pathway by which substances move into cells and so are of great biological and medical importance: a substantial fraction of all drugs used by physicians act on channels. Channels can be studied in the tradition of computational electronics. Drift diffusion equations form an adequate model of IV relations of 6 different channel proteins in ~ 10 solutions over ± 150 mV. Ionic channels can also be studied with the powerful techniques of molecular biology. Atoms can be modified one at a time and the location of every atom can be determined. Ionic channels are natural nanotubes that can be controlled more precisely and easily than physical nanostructures but biologists need help if realistic simulations are to be done at atomic detail.

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Drift-diffusion equations, combined with Poisson's equation (*PDD*), are widely used in physical sciences to describe the flux of charge carriers through systems containing fixed charge (doping) [28]. An important class of biological molecules—proteins called ionic channels—conduct ions (like Na^+ , K^+ , Cl^-), and thus current, through a narrow tunnel of fixed charge ('doping') formed by the polar residues of the protein [36] although electron flow plays no direct role in their conduction of current. These proteins can be studied with the full power of molecular biology [1]; for example, they can be modified one atom at a time

with the techniques of molecular genetics. Thus, these natural nanotubes are a natural 'hole in the wall' that can be controlled more precisely and easily than many physical nanostructures.

Ionic channels open and close ('gate') to give currents that are a random telegraph signal [34]. The properties of gating are complex and the structure(s) and mechanism(s) that produce gating are not known [23], but the flow of ions through open channels is much simpler, and obeys the *PDD* equations, as we shall see [10–13, 19, 37].

Channels are the main pathway by which substances move in and out of cells and so are of

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great biological and medical importance: they are responsible for signaling in the nervous system; for coordination of muscle contraction—including the coordination of cardiac muscle that allows the heart to function as a pump—and they are involved in transport in every cell and organ, for example, in the kidney, intestine and endocrine glands [1, 36]. A substantial fraction of all drugs used by physicians act directly or indirectly on channels [35].

Channels are studied one molecule at a time in hundreds, if not thousands of laboratories every day [14–16, 31], using Neher & Sakmann's patch clamp method [6, 32, 34] (for which they received the Nobel Prize). The concentrations of ions outside channels (that carry current through the channel) can be directly controlled and the shape of current voltage (*IV*) relations can be manipulated. In this way, a wide range of *IV* behavior can be measured from the single doping profile of one type of channel and so much can be inferred about the doping profile from *IV* measurements (if they are taken in many different (pairs of) concentrations of current carriers). For all these reasons, channels are a popular object for experimentation: thousands (!) of abstracts describing their properties are presented each year at the annual meeting of the Biophysical Society and hundreds of papers are published about them, chiefly in the *Biophysical Journal* and the *Journal of Physiology* (London).

Channels are also an appealing and important object for theoretical analysis and numerical simulation. Open channels are probably the simplest protein structures of general biological importance. Unlike many other subjects of biophysical investigation, ionic channels are a general biological system with importance for every organ, tissue, and cell in an animal and plant. Indeed, they are probably just as important for subcellular organelles. Ionic channels are well defined biological systems that can be investigated both with the techniques of molecular biology and of biophysical chemistry, helped substantially by the techniques and insights of semiconductor physics, I believe.

Ionic movement plays an important role in the function of all proteins—e.g., enzymes—and so a model that describes ionic movement in channels is likely to give important insight into protein function in general. Indeed, the closely related Poisson-Boltzmann theory [18, 24] has been of considerable help already, even though it is a strictly equilibrium theory that does not permit flux at any time or location.

Theories of physical chemistry [4] and electrochemistry [30] certainly should be able to predict the movement of ions through a tunnel of fixed charge—a hole in the wall—on the biological time scale of $100\ \mu\text{sec}$ – $10\ \text{sec}$, but they need to be supplemented by the theories and simulations of carrier transport in general, e.g., in semiconductors [28]. The physical chemical tradition has not often dealt with flux [2, 3, 5, 8, 9, 21, 22, 25, 29] in such a system of fixed charge, whereas the flux of charge carriers has been the main subject of semiconductor physics and computational electronics for many years, if not decades.

Five laboratories have shown that the *PDD* of semiconductor physics form an adequate model of *IV* relations of 6 different channel proteins in ~ 10 pairs of solutions (containing different concentrations of the charge carriers Na^+ , K^+ , Cl^-) in the range $\pm 150\ \text{mV}$ [10–13, 37]. The *IV* relations are *qualitatively* different in different types of channels—some are linear, some sublinear and some superlinear—because different channel proteins have qualitatively different profiles of fixed charge arising from their different sequences of amino acids.

The structures of two of these proteins (porin and its mutant G-119D) are known from the standard methods of molecular biology: the location of every atom has been determined by X-ray diffraction [17, 26, 27, 33, 38] with an accuracy of $\sim 0.1\ \text{\AA}$. The mutant has one extra negative charge. Measurements of *IV* relations from a single molecule of porin [37] allow the *PDD* model to estimate the additional charge as $-0.97e$, although this estimate will undoubtedly change as more work is done. (I hasten to add that no information

about the proteins is used in the analysis except the length and diameter of the channel; parameters were not adjusted in any way.)

We conclude that the *PDD* equations seem to be an adequate model open channels. It is surprising that the *PDD* equations work as well as they do, given their evident inadequacies. I imagine they work this well because the fixed charge density of channels is large ($\sim 3 \times 10^{21} \text{ cm}^{-3}$) compared to the concentration of ions outside the channel (2×10^{19} to $1 \times 10^{21} \text{ cm}^{-3}$); because the biological range of voltages is quite limited ($\pm 200 \text{ mV}$); and because the *PDD* model uses effective parameters. Eisenberg, Chen, and Schuss have recently shown how the *PDD* equations can be derived in single file systems like channels that conduct one ion at a time. The *PDD* equations, or equations quite like them, describe the mean properties of ensembles of Langevin equations, each of which specifies the motion of a single ion (of a particular type moving from a given side of the channel), each of which is coupled to its own reaction field described by a Poisson equation and boundary conditions.

The *PDD* equations are just a first, low resolution description of open channels. More realistic models (using Monte Carlo simulations called molecular dynamics in the world of proteins [7, 8, 20]) are needed to provide insight with atomic resolution. It is likely that many critical functions of channels and enzymes will be best understood this way—by atomic resolution simulations that include flux—but biologists cannot do the simulations themselves: their simulations of atomic resolution have been confined strictly to equilibrium as have those of most chemists. Much help is needed from the community of computational electronics if flux, and electrical potentials at the electrodes (i.e., boundaries) of the system, are to be included in chemical simulations of atomic detail.

Ionic channels are so important biologically, but so well defined physically, that they are an ideal object of biophysical investigation. The techniques of computational electronics and molecular biology can be joined together to determine how ionic channels work. Perhaps the same will prove true of

many other chemical and biological systems, but it is wise to try the simple ones first. Nothing is likely to be simpler physically than a hole in the wall.

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Authors' Biography

Bob Eisenberg was born in Brooklyn, NY, in 1942, and educated at Horace Mann School, then Harvard. He was trained as a biochemist by John Edsall, as a biophysicist by Bernard Katz, as a physiologist by Andrew Huxley, and (from afar) by Alan Hodgkin. Working at UCLA from 1968–1976, he learned some applied mathematics from Julian Cole and Victor Barcilon while studying current flow in cells, chiefly muscle. Moving to Rush Medical College, Chicago, in 1976, he became the Bard Professor and Chairman of their Department of Molecular Biophysics and Physiology. He has been working on self-consistent models of ionic channels for some ten years.