

Membranes in Optic Nerve Models

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Abstract

Membranes are present in all cells and tissues. Mathematical models of cells and tissues need a compact mathematical description of membranes with a resolution of about 1 nm. Membranes isolate cells because ions have difficulty penetrating the dielectric barrier they create. Here we introduce a dielectric mathematical membrane condition to replace a condition that did not include dielectric properties. Our mathematical membrane condition includes a dielectric lipid bilayer punctured by channels that conduct ions selectively.

1 Introduction

Membranes isolate cells and organelles from the solutions around them. Membranes are lipid bilayers that isolate regions because of their polarization. The energy to move an ion from salt solution into the lipid of the bilayer is very large. As it moves, an ion must be separated from its ionic atmosphere (measured by the Debye length) and dielectric surround (measured by the Bjerrum

length) that balance its permanent charge. The dielectric lipid has zero ionic atmosphere and little dielectric atmosphere to balance its charge. The movement of the ion from solution into the lipid of the membrane involves a huge energy barrier. The energy barrier isolates the interior of cells and organelles from their surround. Membranes are punctured by channels that provide pathways for ion flow with relatively small energy barriers (when the channel is open). In recent papers on the optic nerve [1], we followed Mori [2, 3] and described the membranes by a condition involving just the channels and their conductance

$$\mathbf{J}_k^i \cdot \mathbf{n} = \frac{g_m^i}{z^i e} (\phi_{in} - \phi_{ex} - E^i), \quad (1)$$

where $i = \text{Na}^+, \text{K}^+, \text{Cl}^-$, $E^i = \frac{k_B T}{z^i e} \log \left(\frac{c_{ex}^i}{c_{in}^i} \right)$ is the Nernst potential.

Here we introduce a mathematical membrane condition involving membrane dielectric properties (represented realistically as a capacitance) as well as membrane conductance. We show that our previous results are not significantly changed by introducing the new mathematical membrane condition. We choose the resolution of the membrane condition so it describes membrane properties important in Hodgkin Huxley computations of the action potential and calculations of the associated ionic and water fluxes. Other effects, such as the ionic atmosphere surrounding the membrane, or the permanent charge inside the ion channel require higher resolution models. We will use the realistic membrane condition (that includes capacitance) in future work lest readers and students misunderstand the physical basis of the main function of membranes, namely to provide a large dielectric energy barrier that isolates cells and organelles from their surrounds.

2 Mathematical membrane condition

In this section, we discuss the effect of the new mathematical membrane condition for both axon cell membrane and glial cell membrane.

The cell membrane can be viewed as a capacitor representing the lipid dielectric and a set of conductances representing ion channels following ample precedent. In the model, we introduce the $J_{m,k}^i$, ($k = gl, ax, i = \text{Na}^+, \text{K}^+, \text{Cl}^-$) for this ion flux and the transmembrane source term in the [1] can be written as

$$J_k^{m,i} = J_{p,k}^i + J_{c,k}^i + J_{m,k}^i, \quad k = gl, ax, \quad i = \text{Na}^+, \text{K}^+, \text{Cl}^-.$$

where

$$J_{m,k}^i = \lambda^i \frac{C_m}{z^i e} \frac{d(\phi_k - \phi_{ex})}{dt},$$

and active ion pump source $J_{p,k}^i$ and passive ion channel source $J_{c,k}^i$, on the k th membrane. The conservation of ion concentration implies the following system of partial differential equations to describe the dynamics of ions in each region,

for $i = \text{Na}^+, \text{K}^+, \text{Cl}^-$

$$\begin{aligned} \frac{\partial (\eta_{gl} c_{gl}^i)}{\partial t} + \mathcal{M}_{gl} J_{gl}^{m,i} + \nabla \cdot (\eta_{gl} \mathbf{j}_{gl}^i) &= 0, & \text{in } \Omega_{OP}, \\ \frac{\partial (\eta_{ax} c_{ax}^i)}{\partial t} + \mathcal{M}_{ax} J_{ax}^{m,i} + \frac{\partial}{\partial z} (\eta_{ax} j_{ax,z}^i) &= 0, & \text{in } \Omega_{OP}, \\ \frac{\partial (\eta_{ex} c_{ex}^i)}{\partial t} - \mathcal{M}_{ax} J_{ax}^{m,i} - \mathcal{M}_{gl} J_{gl}^{m,i} + \nabla \cdot (\eta_{ex} \mathbf{j}_{ex}^i) &= 0, & \text{in } \Omega_{OP}, \end{aligned}$$

The electric potentials in the axon compartment, glial compartment and extracellular space are

$$\begin{aligned} \sum_i z^i e \mathcal{M}_{gl} J_{gl}^{m,i} + \sum_i z^i e \nabla \cdot (\eta_{gl} \mathbf{j}_{gl}^i) &= 0, \\ \sum_i z^i e \mathcal{M}_{ax} J_{ax}^{m,i} + \sum_i z^i e \frac{\partial}{\partial z} (\eta_{ax} j_{ax,z}^i) &= 0, \\ \sum_i z^i e \nabla \cdot (\eta_{gl} \mathbf{j}_{gl}^i) + \sum_i z^i e \frac{\partial}{\partial z} (\eta_{ax} j_{ax,z}^i) + \sum_i z^i e \nabla \cdot (\eta_{ex} \mathbf{j}_{ex}^i) &= 0, \end{aligned}$$

3 Simulation

Here we repeat the simulations of the optic nerve published previously but with modified parameters. We choose the capacitance on both axon and glial membrane as $C_m = 7.5 \times 10^{-3} \text{ F} \cdot \text{m}^{-2}$ [4] and $\lambda^i = \frac{1}{3}$ for each ion specie. Furthermore, action potentials are generated by the axon with the parameters in the following Table 1.

Parameters	Previous Value	New Value
\mathcal{M}_{ax}	$5.98 \times 10^6 \text{ m}^{-1}$	$2.392 \times 10^5 \text{ m}^{-1}$
$I_{ax,1}$	$9.56 \times 10^{-4} \text{ A/m}^2$	$2.39 \times 10^{-2} \text{ A/m}^2$
$I_{ax,2}$	$1.3 \times 10^{-4} \text{ A/m}^2$	$3.25 \times 10^{-3} \text{ A/m}^2$
g_{leak}^{Na}	$4.8 \times 10^{-3} \text{ S/m}^2$	$1.2 \times 10^{-1} \text{ S/m}^2$
g_{leak}^K	$2.2 \times 10^{-2} \text{ S/m}^2$	$5.5 \times 10^{-1} \text{ S/m}^2$
\bar{g}^{Na}	$1.357 \times 10^1 \text{ S/m}^2$	$3.393 \times 10^2 \text{ S/m}^2$
\bar{g}^K	2.945 S/m^2	$7.364 \times 10^1 \text{ S/m}^2$
g_{ax}^{Cl}	$1.5 \times 10^{-1} \text{ S/m}^2$	3.75 S/m^2
I_{shock}	$3 \times 10^{-3} \text{ A} \cdot \text{m}^{-2}$	$7.5 \times 10^{-2} \text{ A} \cdot \text{m}^{-2}$

Table 1: Parameters Table

Remark 3.1. *The parameters are adjusted based on the following facts*

1. *The size of necturus optic nerve is much smaller that squid;*

2. In [1], we estimated the membrane conductance based on the experimental results [5] and the fact that the density of ion channel in *Necturus* ($24/\mu\text{m}^2$) is lower than that of squid ($170 - 290/\mu\text{m}^2$). Without the capacitance effect, the conductance of axon may be underestimated.

At the same time, in [5], the measured K^+ variation is around 0.2mM for single action potential. In order to match this, we keep the value $\mathcal{M}_{ax} * \bar{g}_k$ to be the same, while decrease the value of \mathcal{M}_{ax} and increase the value of conductance. The other parameters are adjusted based on the rest state ion concentrations.

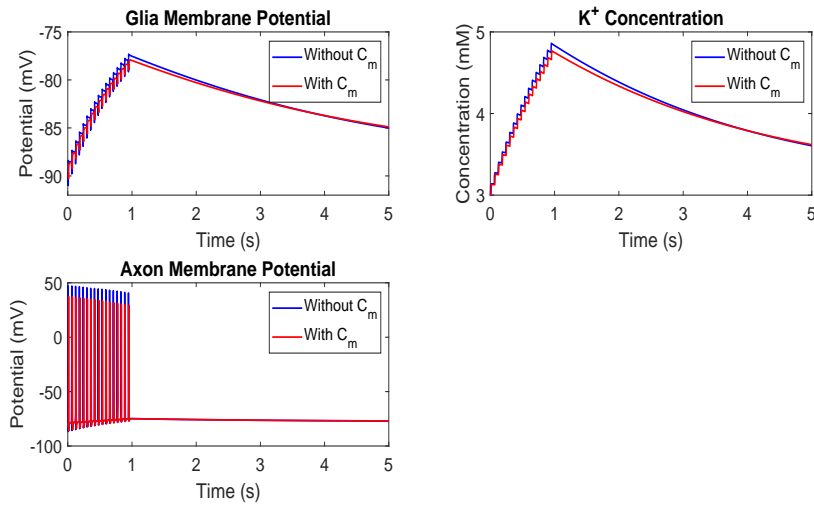


Figure 1: Comparison between the model with and without cell membrane capacitance effect.

Fig. 1 shows results from simulations using the parameters in Table 1. We simulate Orkand's experimental setup with the bath solution containing 3mM K^+ . In Figure 1, we compare the glial membrane potential, potassium concentration in the extracellular space and axon membrane potential. The results do not differ significantly from those published using Mori's purely conductive model of the membrane.

References

- [1] Y. Zhu, S. Xu, R. S. Eisenberg, H. Huang, Optic nerve microcirculation: Fluid flow and electro-diffusion, arXiv preprint arXiv:2101.12152 (2021).
- [2] Y. Mori, C. Liu, R. S. Eisenberg, A model of electrodiffusion and osmotic water flow and its energetic structure, *Physica D: Nonlinear Phenomena* 240 (22) (2011) 1835–1852.

- [3] Y. Mori, A multidomain model for ionic electrodiffusion and osmosis with an application to cortical spreading depression, *Physica D: Nonlinear Phenomena* 308 (2015) 94–108.
- [4] G. Major, A. U. Larkman, P. Jonas, B. Sakmann, J. . Jack, Detailed passive cable models of whole-cell recorded ca3 pyramidal neurons in rat hippocampal slices, *Journal of Neuroscience* 14 (8) (1994) 4613–4638.
- [5] R. Orkand, J. Nicholls, S. Kuffler, Effect of nerve impulses on the membrane potential of glial cells in the central nervous system of amphibia., *Journal of neurophysiology* 29 (4) (1966)